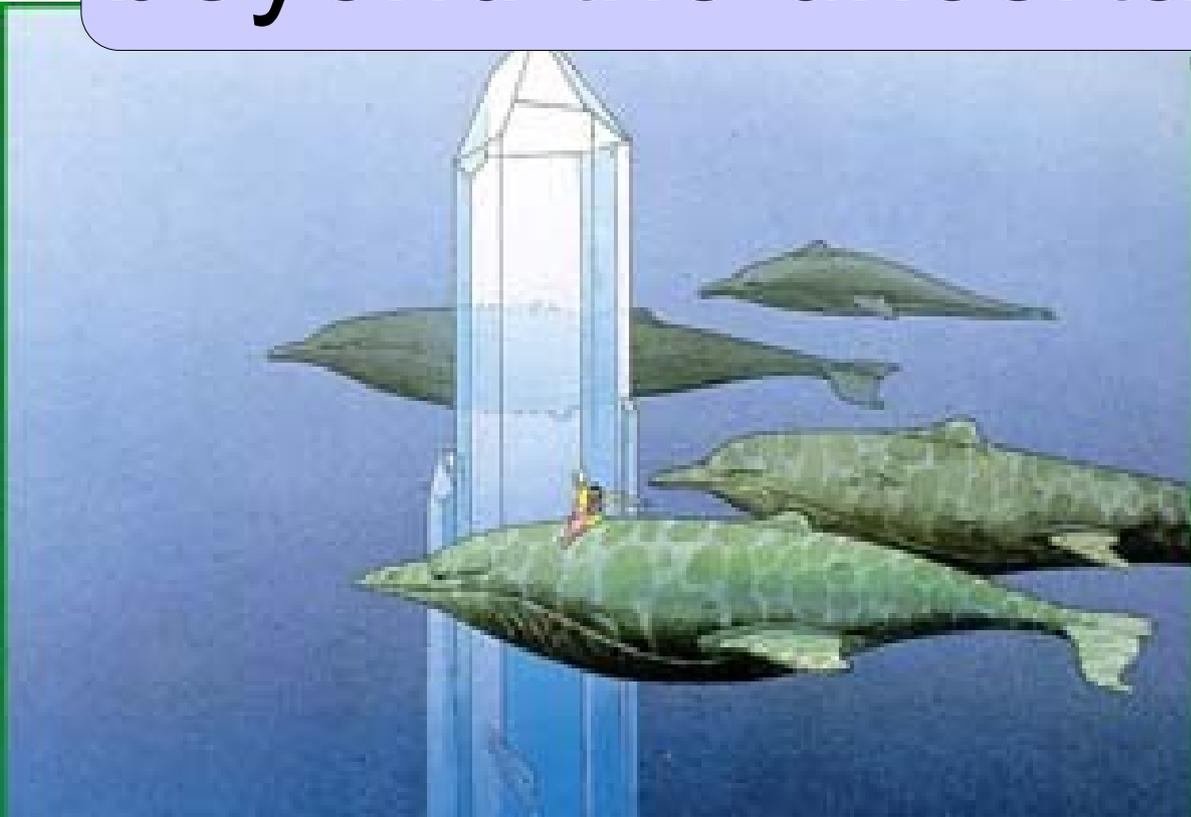


# Quantum measurement bounds: beyond the uncertainty relations



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quantum information  
theory group  
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What am I talking about?



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**A new measurement-precision  
bound:**

**a generalization of Heisenberg-type  
uncertainty relations**

# Outline

- Heisenberg uncertainty relations and q. Cramer-Rao bounds
- Our bound: **a first moment generalization**
- Ideas behind the proof
- Prior information?

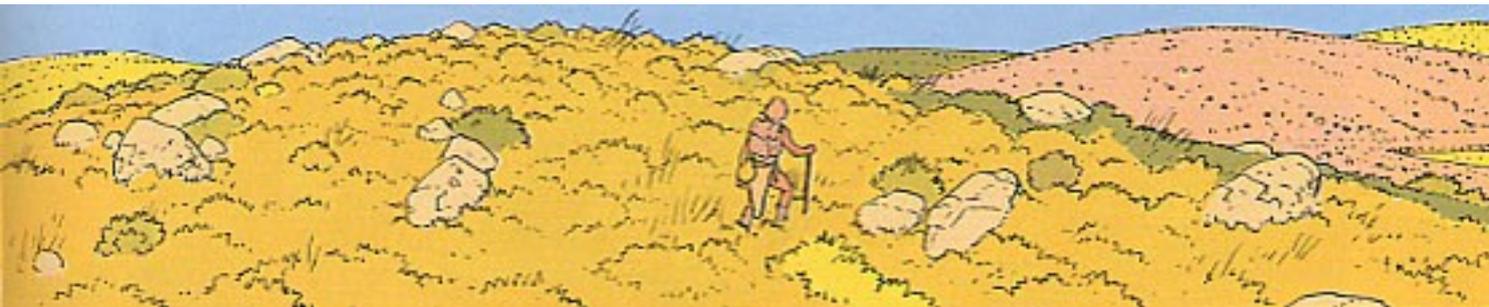
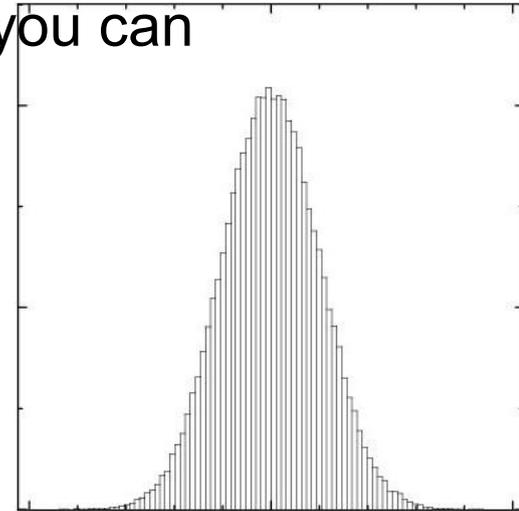


**Abstract:** In quantum mechanics, the Heisenberg uncertainty relations and the Cramer-Rao inequalities typically limit the precision in the estimation of a parameter through the *standard deviation* of a conjugate observable. Here we extend these relations by giving a bound to the precision of a parameter in terms of the *average value* of the conjugate observable. This has both foundational and practical consequences: in quantum optics it resolves a controversy on which is the ultimate precision in interferometry.

## Heisenberg uncertainty relation

“If you have a probe system with spread  $\Delta p$  in momentum, you can discover its position with uncertainty  $\Delta x$ ”

$$\Delta X \geq \frac{\hbar}{2\Delta p} \equiv \frac{1}{\Delta H}$$

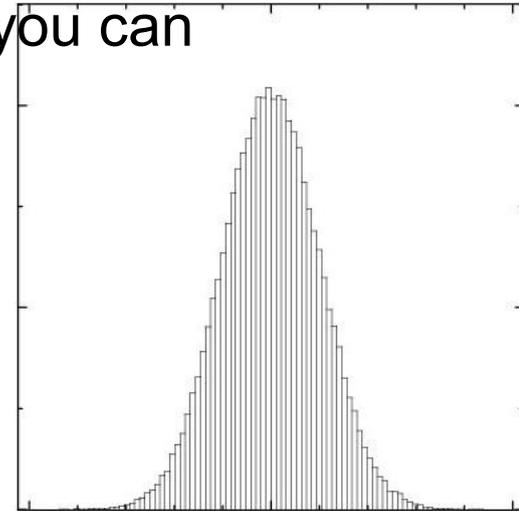


# Heisenberg uncertainty

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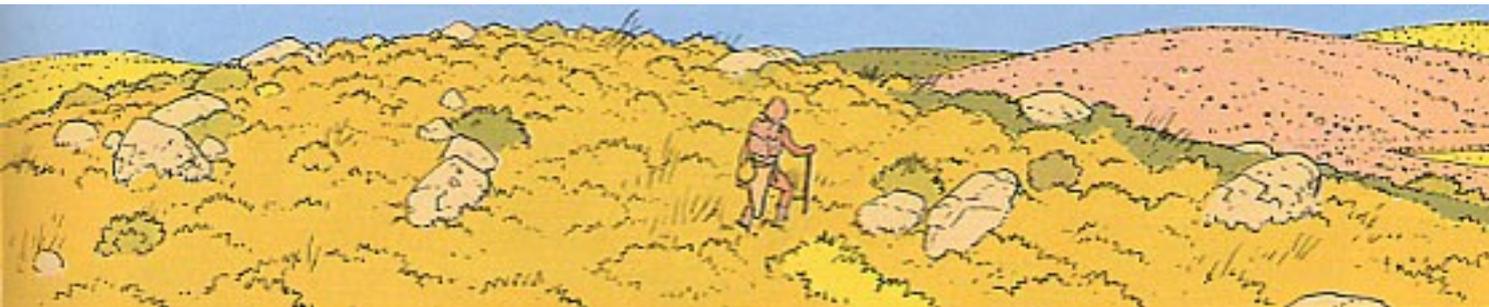
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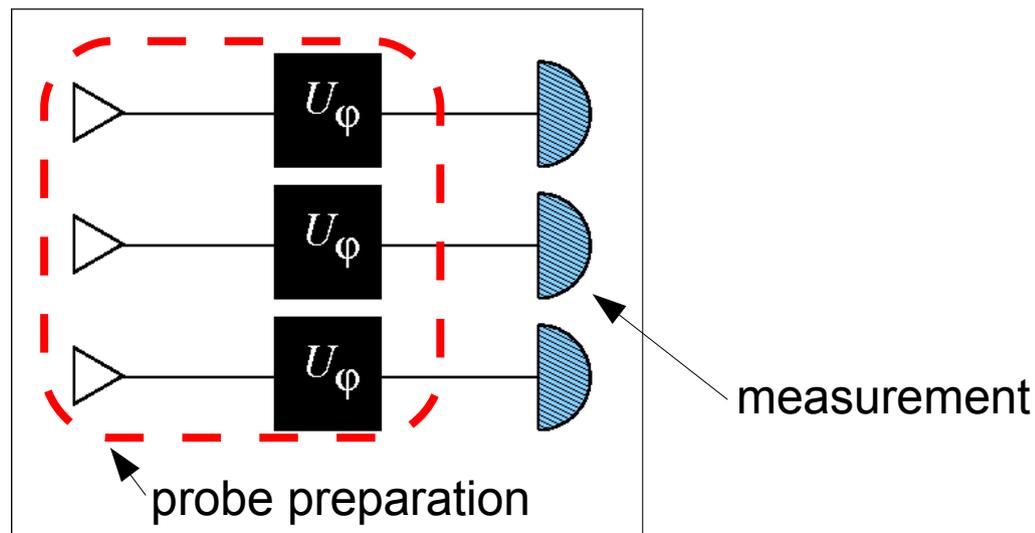
$H$  is the generator of translations of  $X$ :

$$\rho_x = e^{-ixH} \rho_0 e^{ixH}$$



# Heisenberg uncertainty

To increase precision, prepare and repeat the measurement  $\nu$  times,

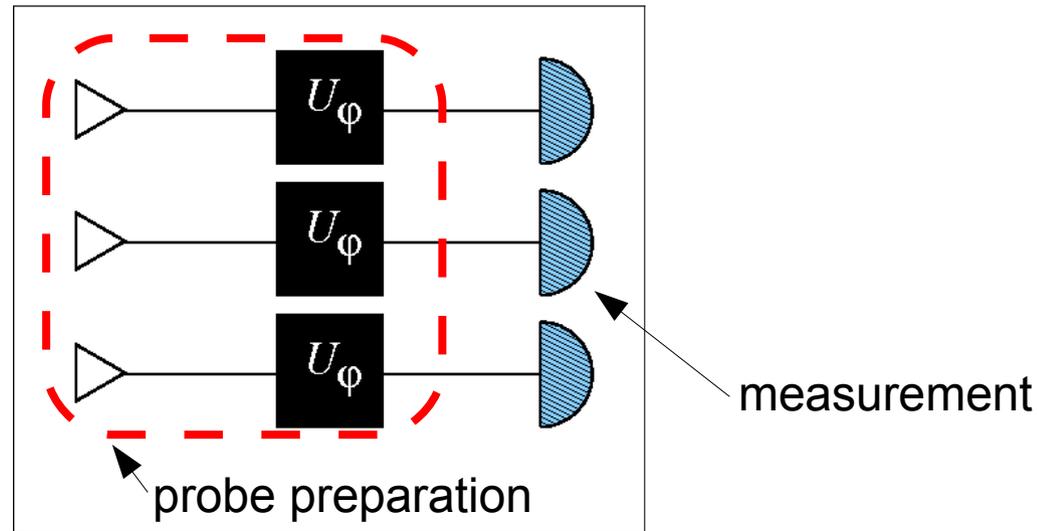


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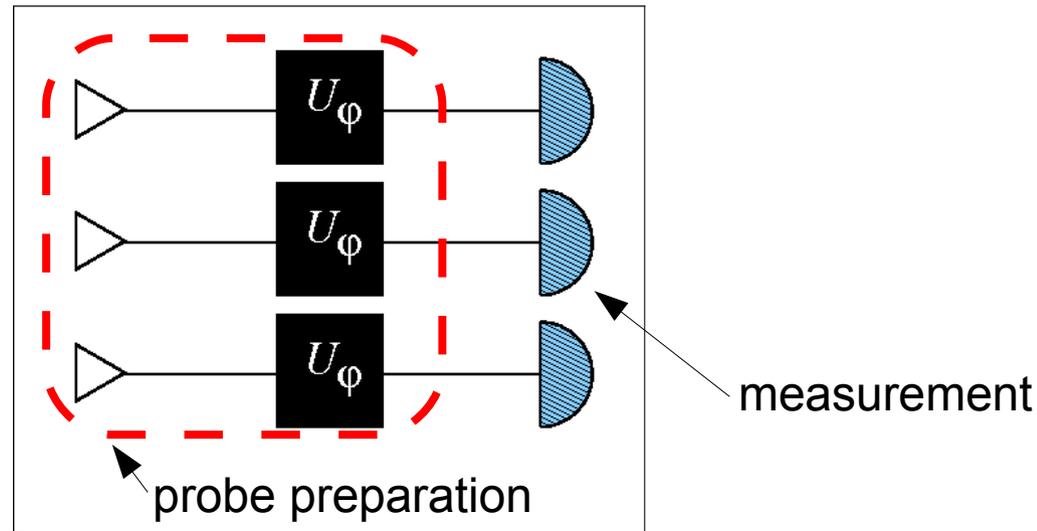
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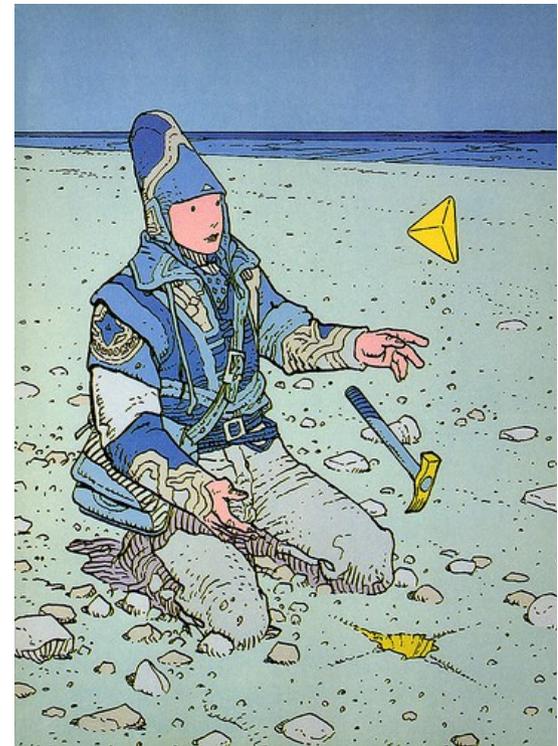
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Remember, we're looking for the **AVERAGE** position (not *the* position)

# Cramer-Rao bound

An **achievable** lower bound on the precision for the measurement of a **parameter**



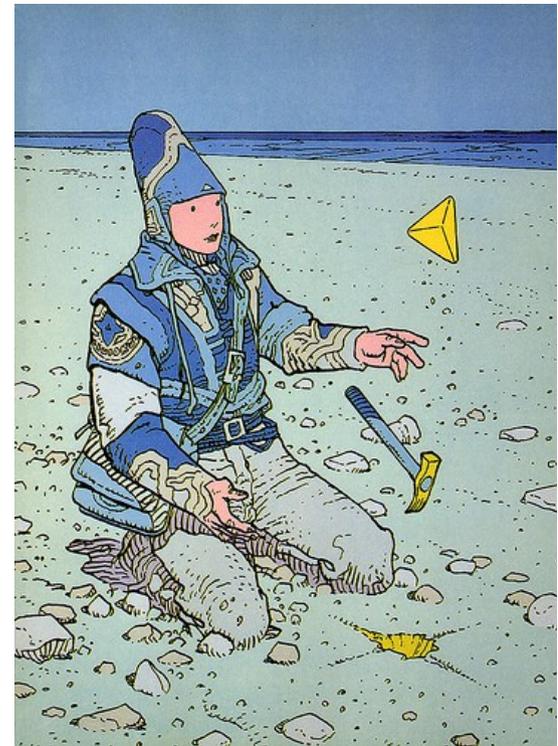
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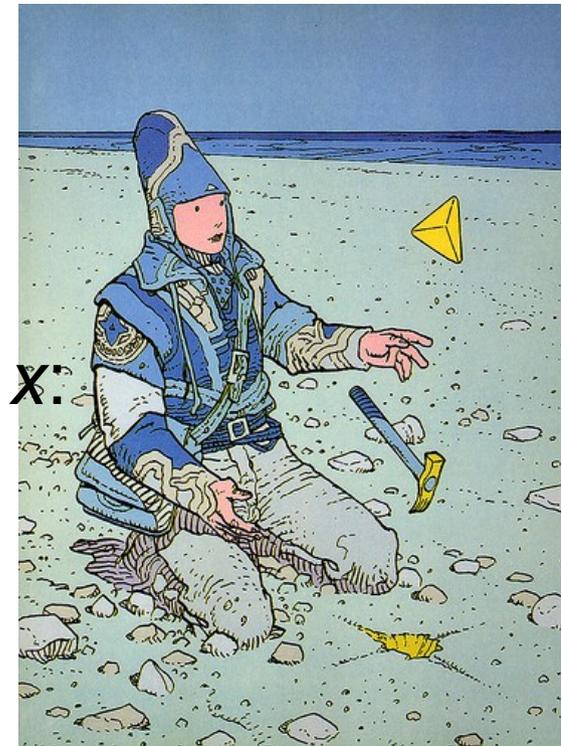
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but now  $x$  is a **parameter** (not an observable)

and again  $H$  is the generator of translations of  $x$ :

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# New bound

Heisenberg / Cramer-Rao

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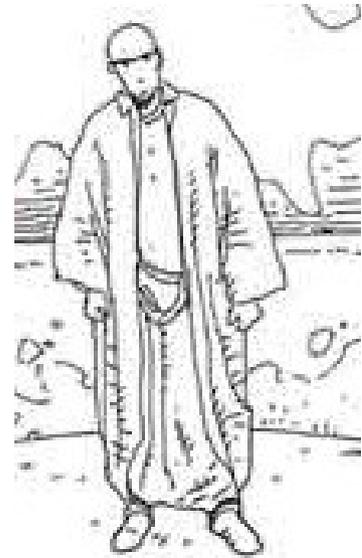
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ground state (minimum eigenvalue of  $H$ )



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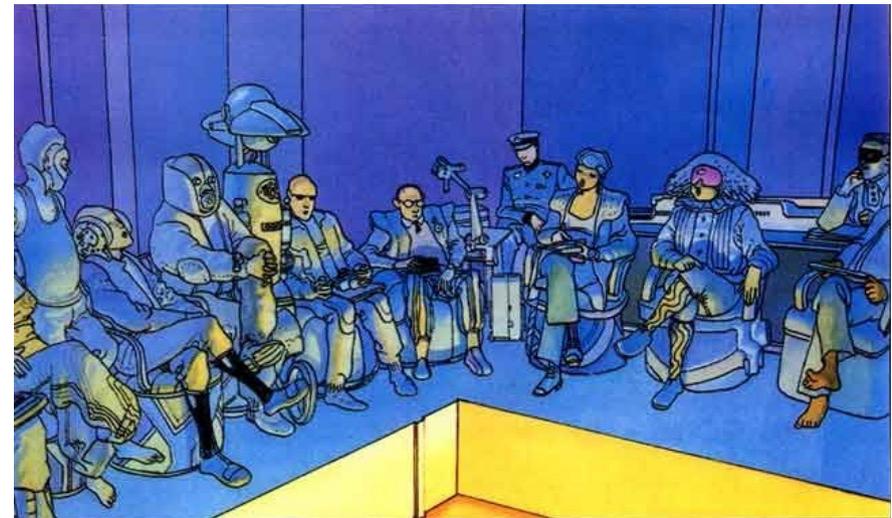


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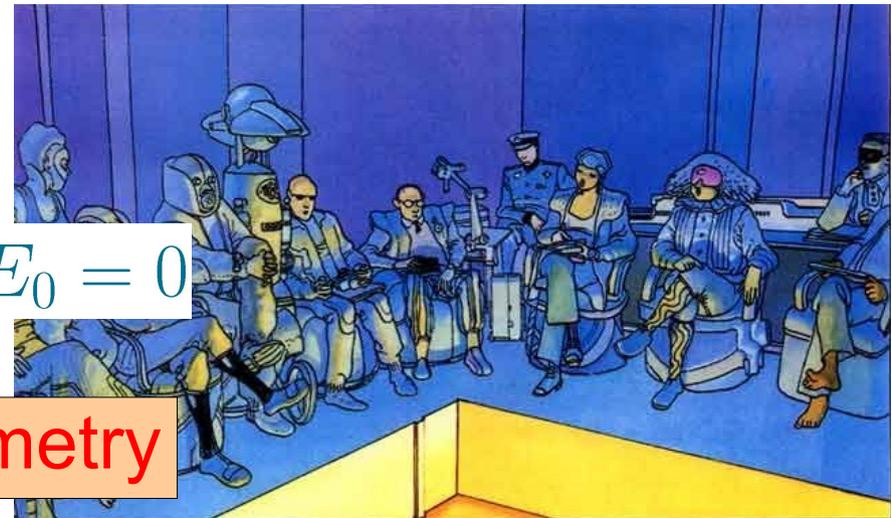
**a first moment generalization of  
Heisenberg uncert. / Cramer-Rao bound**

For interferometry:

$$H = a^\dagger a, \quad x = \phi, \quad \nu = 1$$

$$\Delta \phi \geq \frac{\kappa}{N}$$

$$N = \langle a^\dagger a \rangle, \quad E_0 = 0$$



**The Heisenberg bound for interferometry**

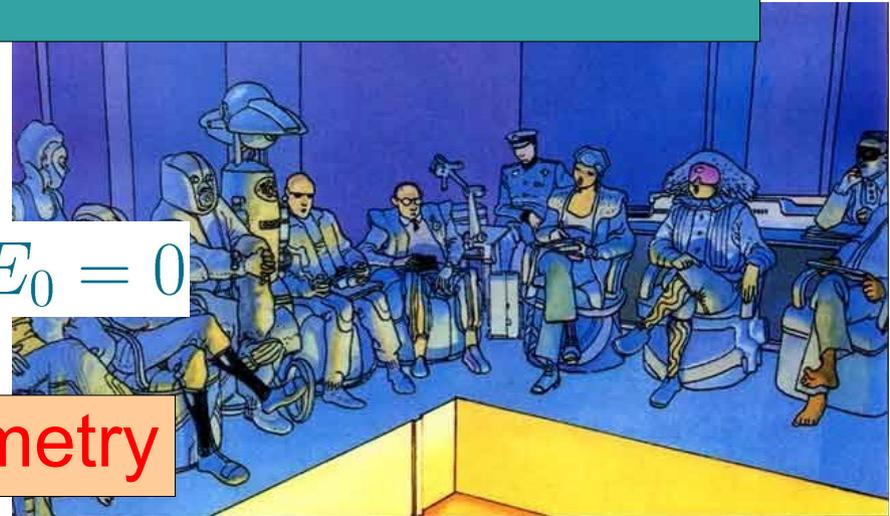
People are saying the Heisenberg bound **CAN** be beaten...

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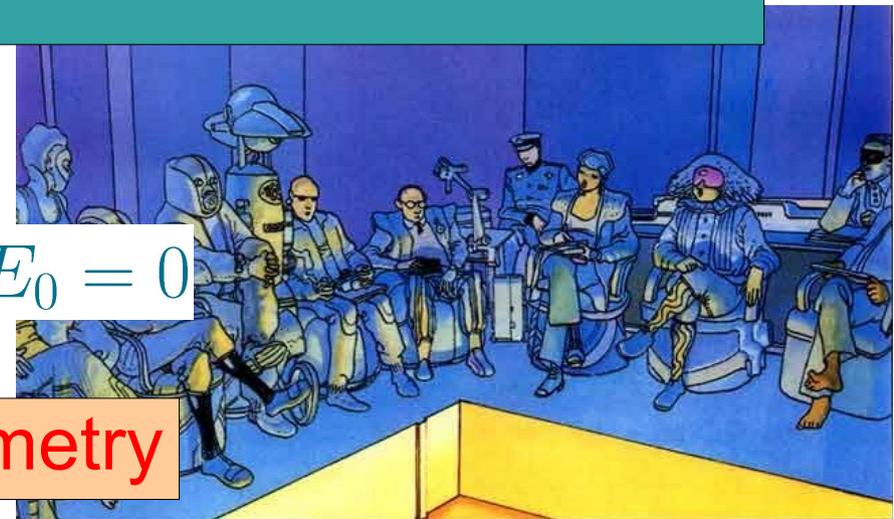
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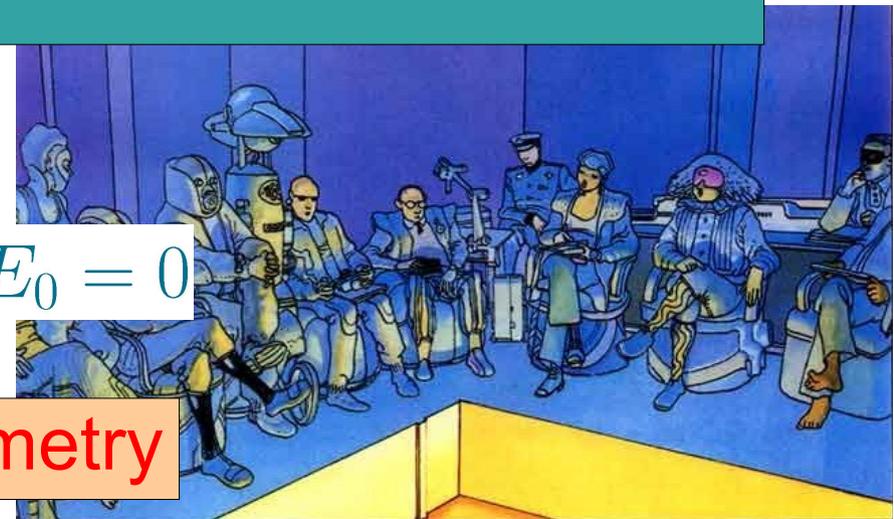
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**BOTH!!!**

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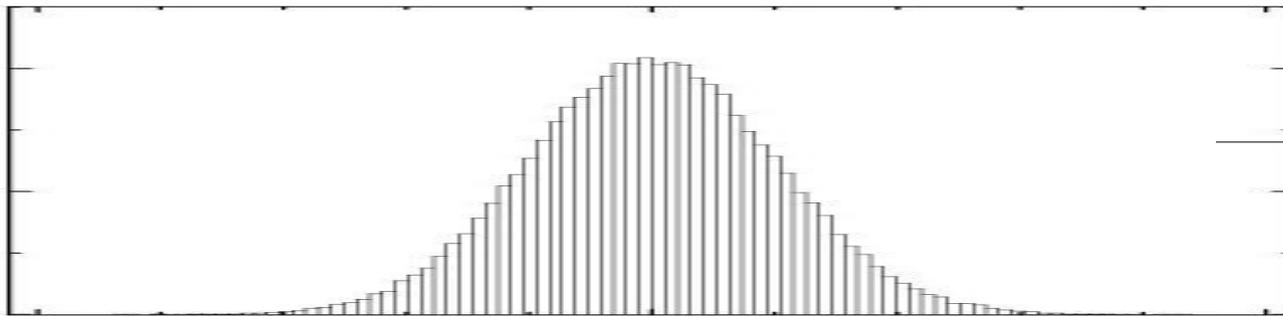
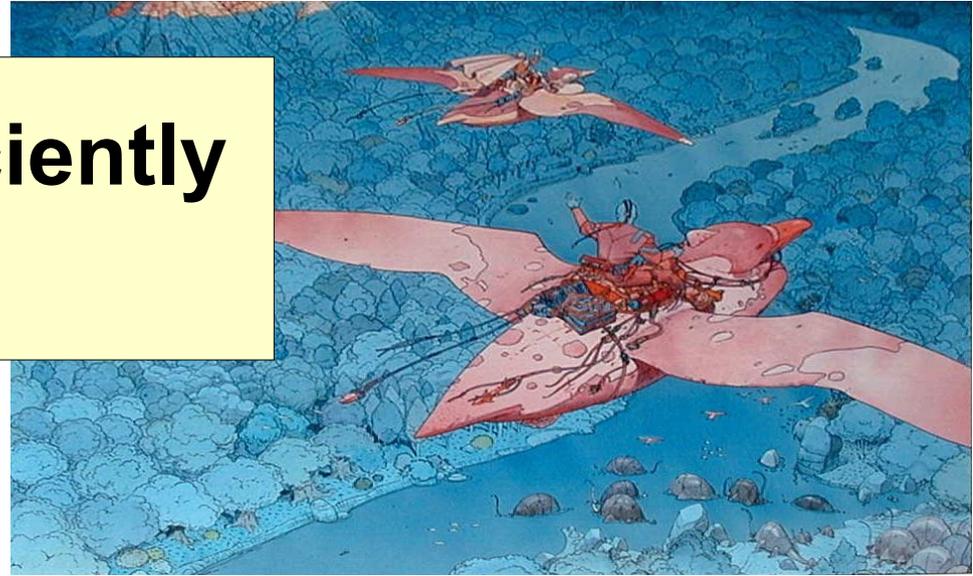
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# Hypothesis behind our bound

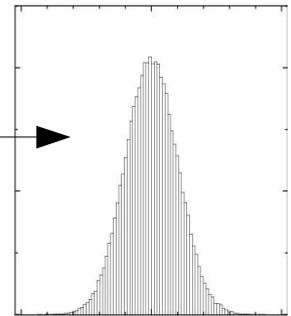
Our bound works for **sufficiently good** estimation strategies

i.e. that give much better precision than the prior information:

(I'll get back to this later)



3x to 4x



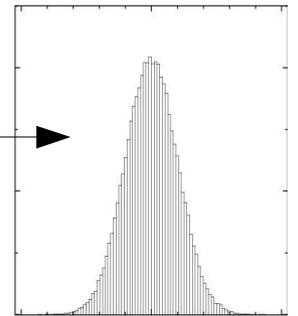
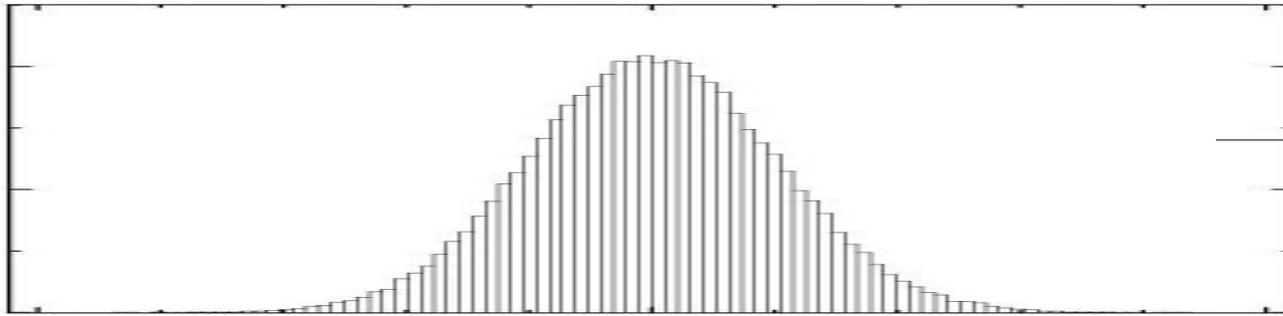
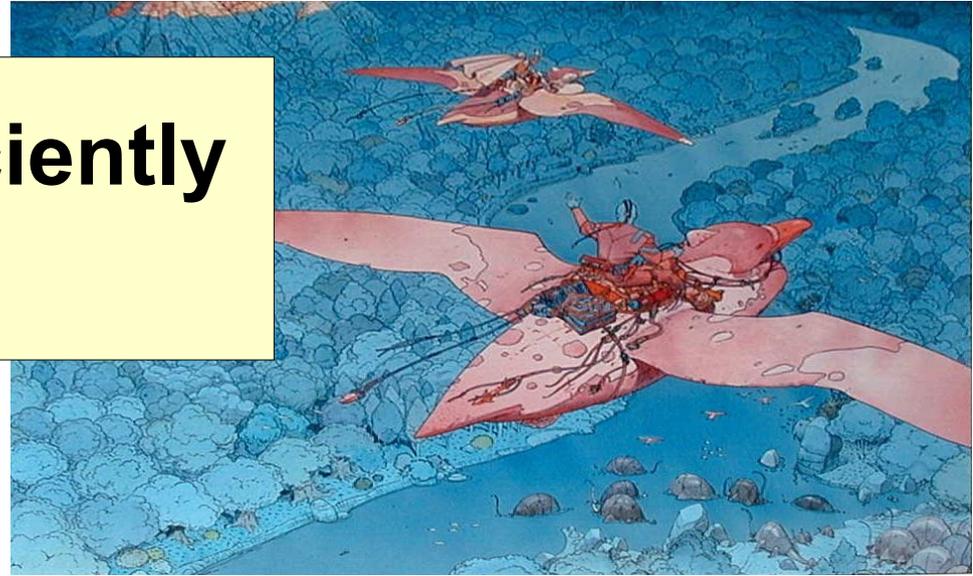
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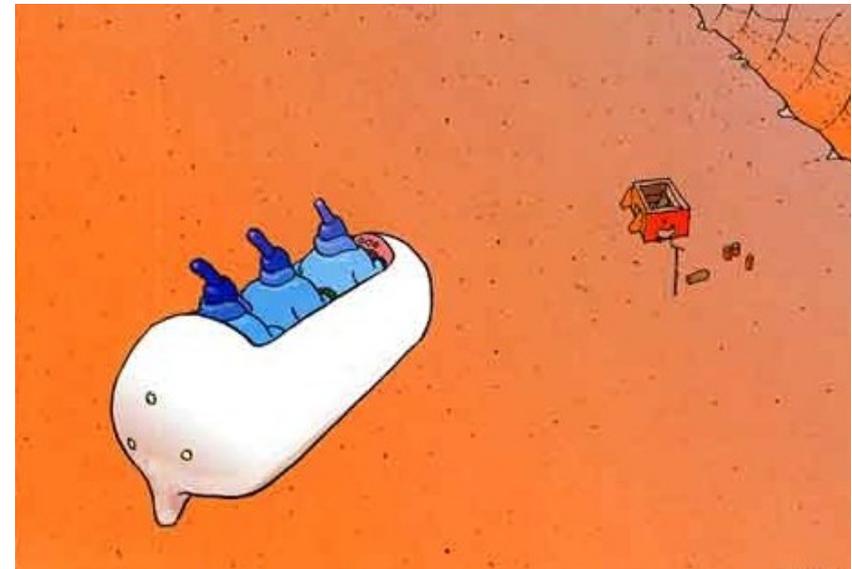


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→ **UNINTERESTING** strategies

# Parameter-dependent error

..but  $\Delta X$  may depend on the (true value of the) parameter  $x$   
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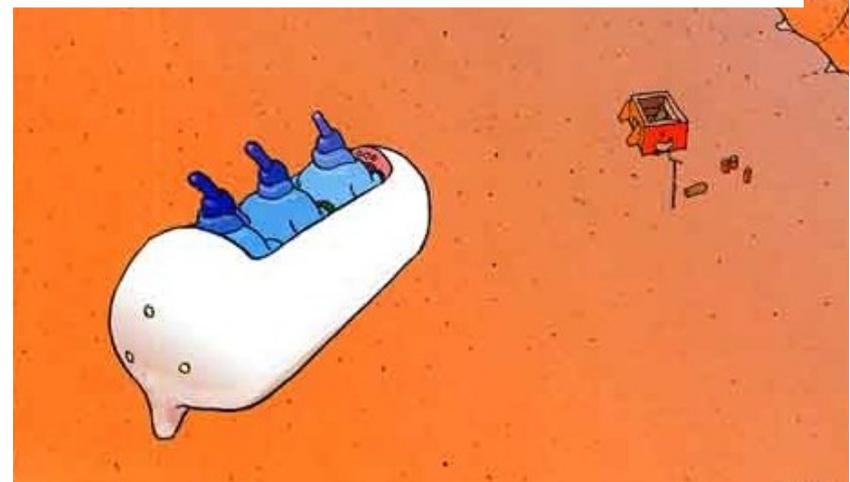
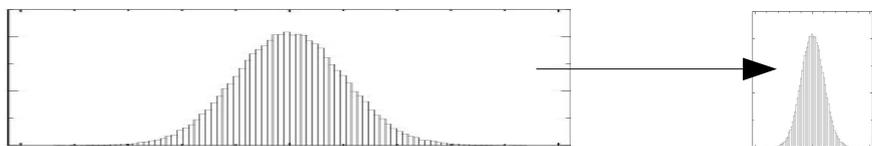
..but  $\Delta X$  may depend on the (true value of the) parameter  $x$   
(so the error can be zero for a specific value of  $x$ ; e.g. a broken clock!)

$\Rightarrow$  our bound holds for the **average**

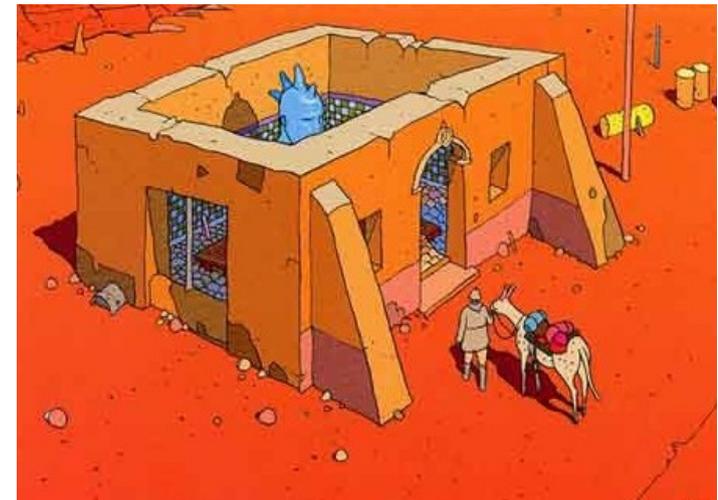
$$\frac{\Delta X(x) + \Delta X(x')}{2} \geq \frac{\kappa}{\sqrt{\nu \mathcal{H}}}$$

**IF** we can find  $x$  and  $x'$  sufficiently far apart

we can always do that if the estimation strategy is good enough, i.e. it yields much better precision than the prior information:

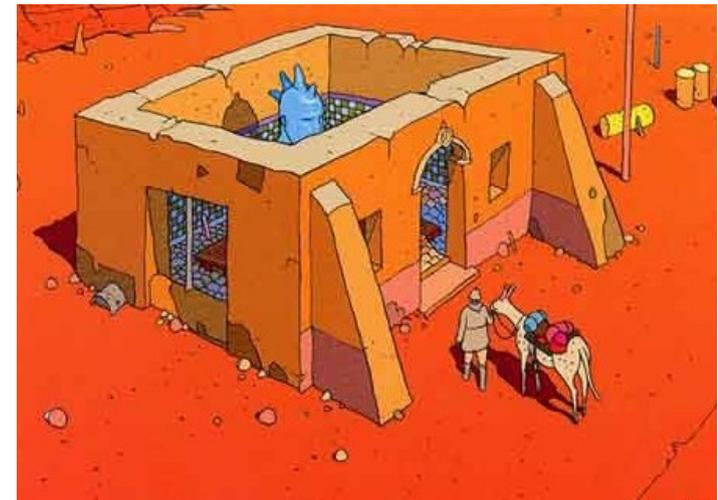


The bound holds generally:



## The bound holds generally:

- Unbiased or biased measurement procedures
- Local or global estimation strategies
- Pure or mixed state of the probes.



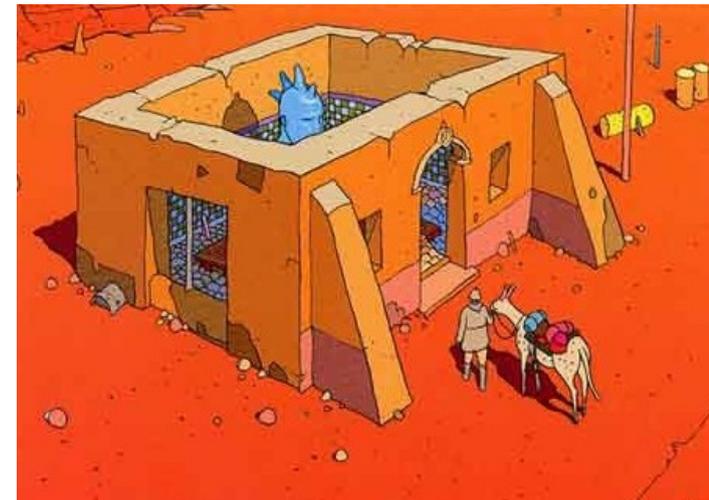
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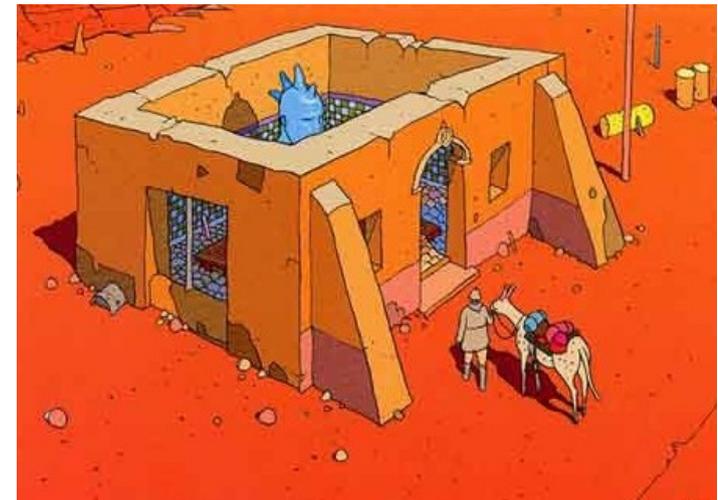
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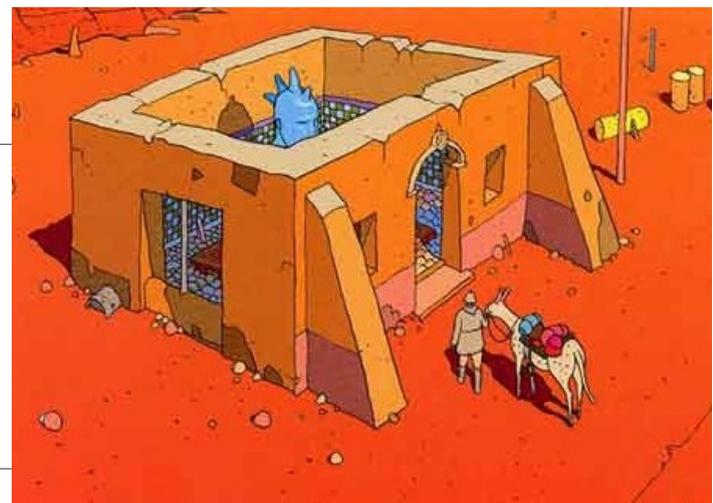
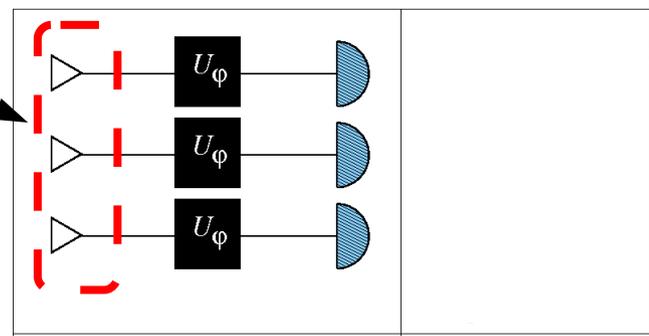
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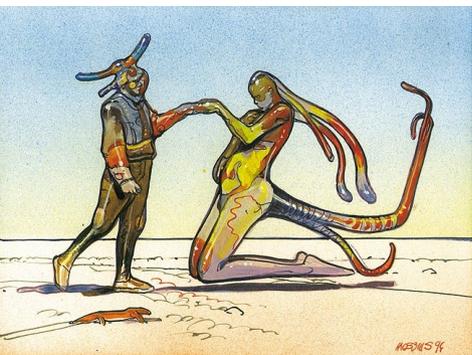
$$\Delta^2 X = \sum_y p(y) [y - x]^2$$

$$p(y) = \left| \langle y | [\psi(x)]^{\otimes \nu} \rangle \right|^2 \leftarrow \text{prob. of obtaining } y \text{ when the true value of the } \nu \text{ probes was } x$$



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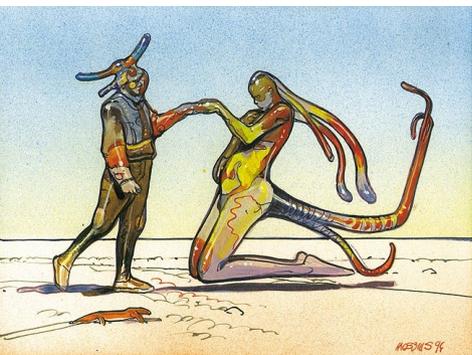
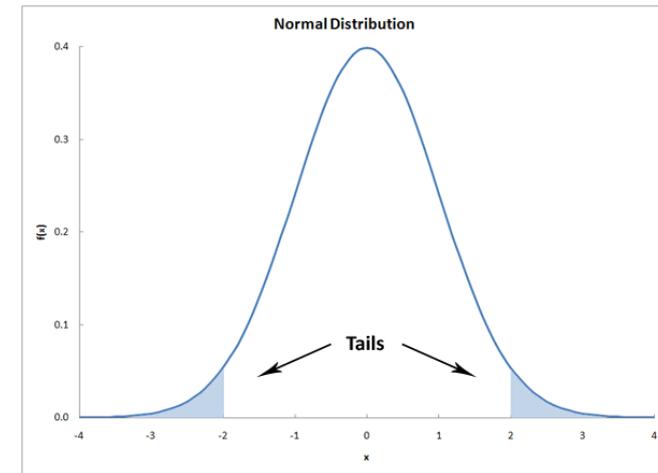
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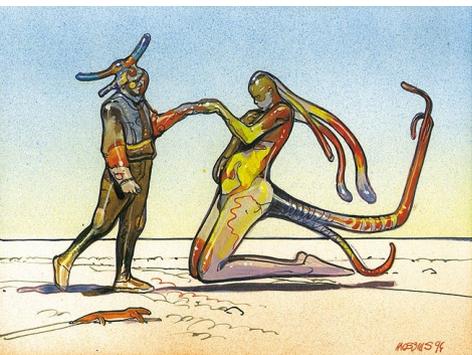
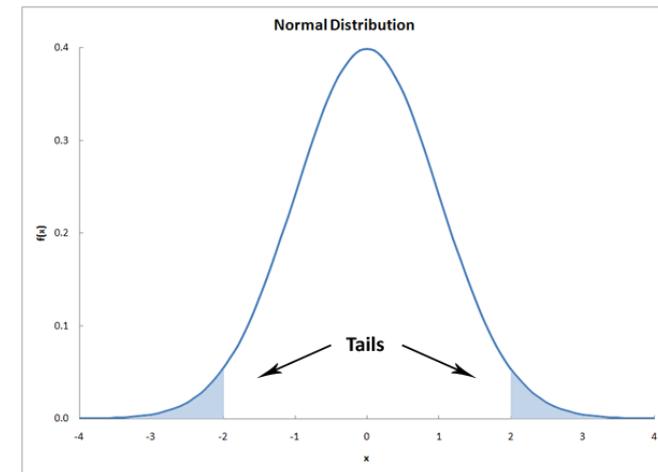


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$\Rightarrow$  the prob that outcome lies within  $\lambda\Delta X$  of the mean value of  $|\psi(x)\rangle^{\otimes \nu}$  is  $\leq 1/\lambda^2$

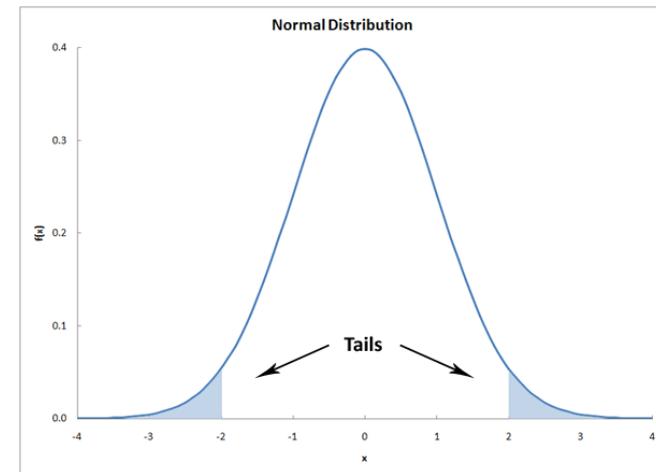


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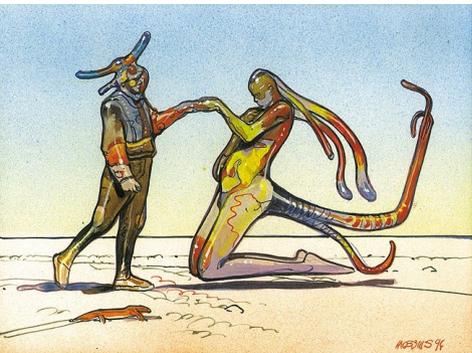
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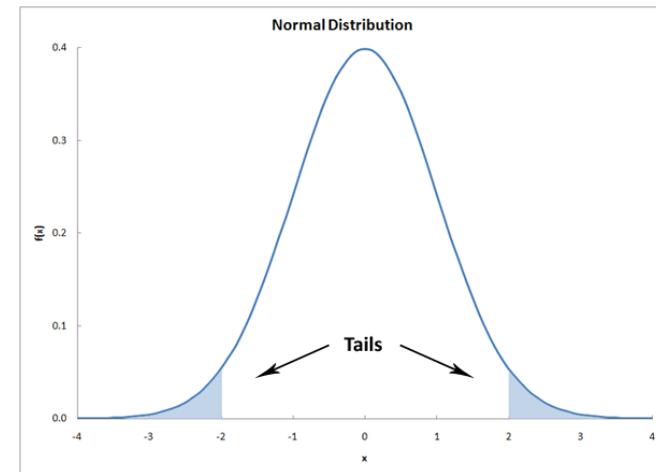


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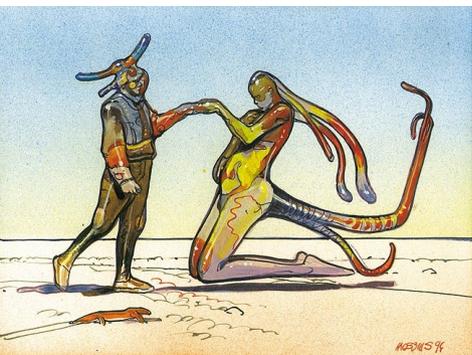
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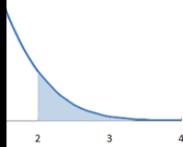
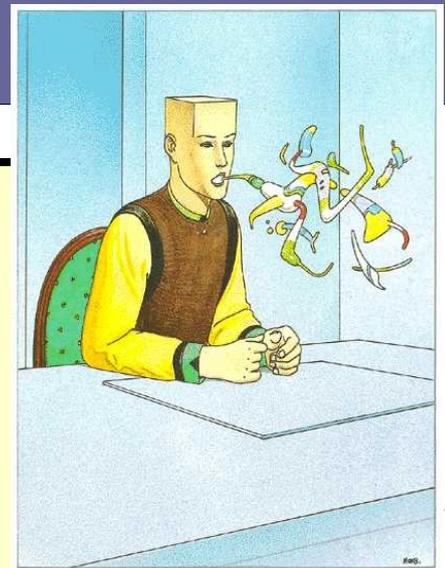
the prob that outcome lies within  $\lambda\Delta X$  of the mean value of  $|\psi(x')\rangle^{\otimes \nu}$  is  $\leq 1/\lambda^2$

⇒ the overlap cannot be too large!

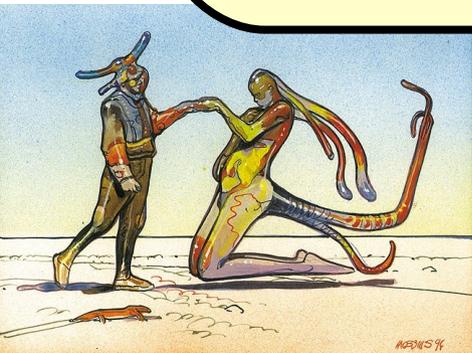


$$F = |\langle \psi(x) | \psi(x') \rangle|^{2\nu} \leq \frac{4}{\lambda^2}$$

LOST?



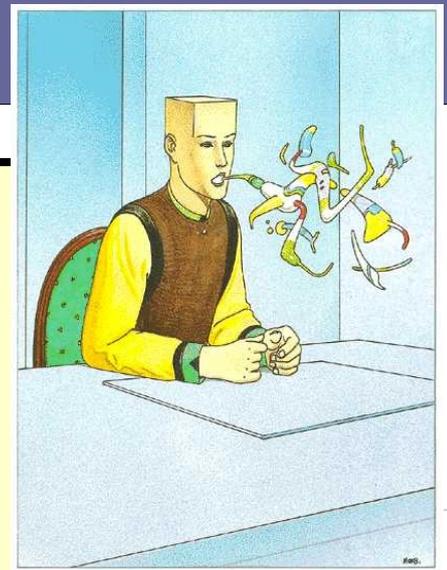
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C

Tch





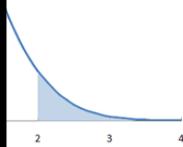
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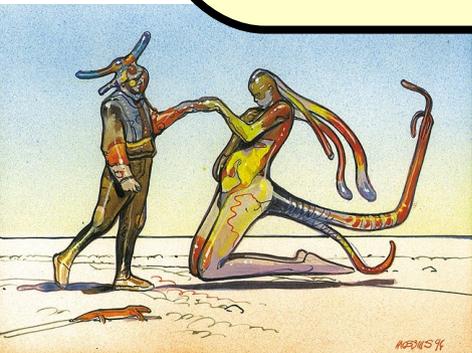
connect the error  $\Delta X$  to the fidelity  $F$  :

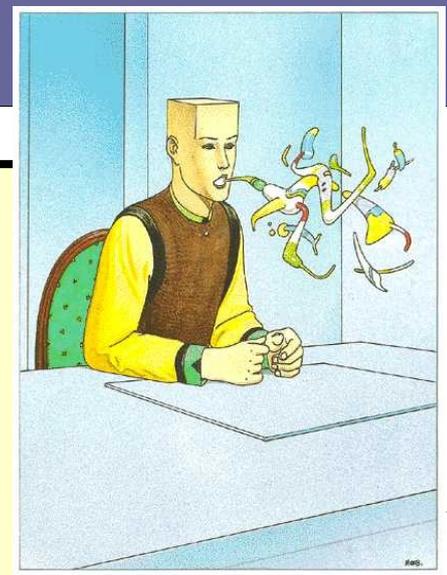
$$x' - x = 2\lambda\Delta X \Rightarrow$$

$$F \leq \frac{4}{\lambda^2}$$



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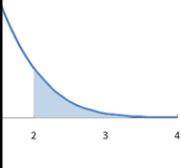
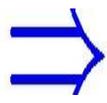
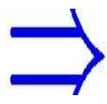
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now, we have to connect the fidelity with the expectation value  $\mathcal{H}$  of the translation generator

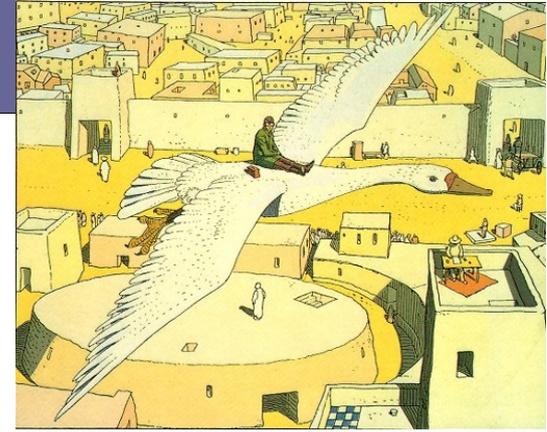
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C  
Tch



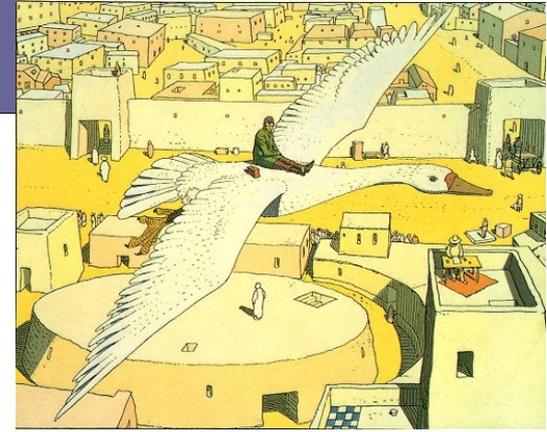
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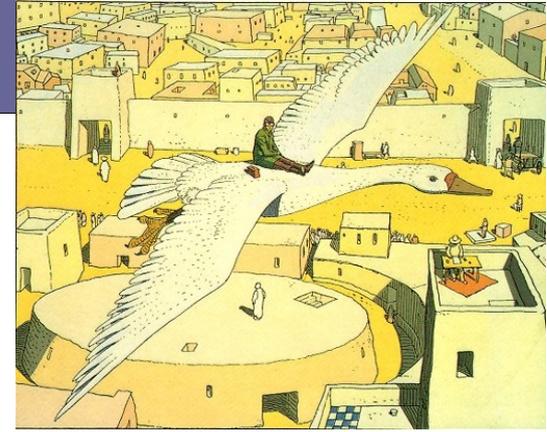
Use the q. speed limit



Use the q. speed limit

Consider an “evolution”  $|\psi(x')\rangle = e^{-i(x'-x)H} |\psi(x)\rangle$





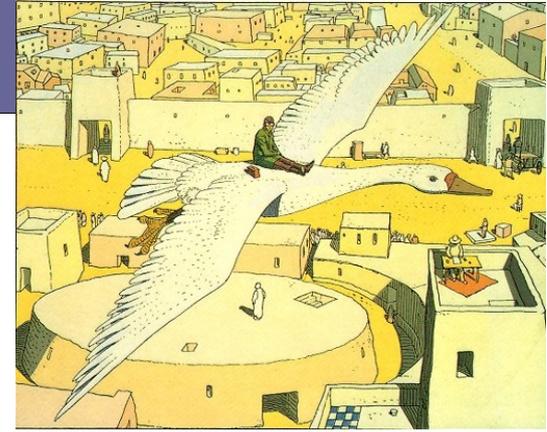
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Consider an “evolution”  $|\psi(x')\rangle = e^{-i(x'-x)H} |\psi(x)\rangle$ , what is the smallest “time”  $x'$  at which the fidelity

$$F \equiv |\langle \psi(x) | \psi(x') \rangle|^2 = \epsilon ?$$

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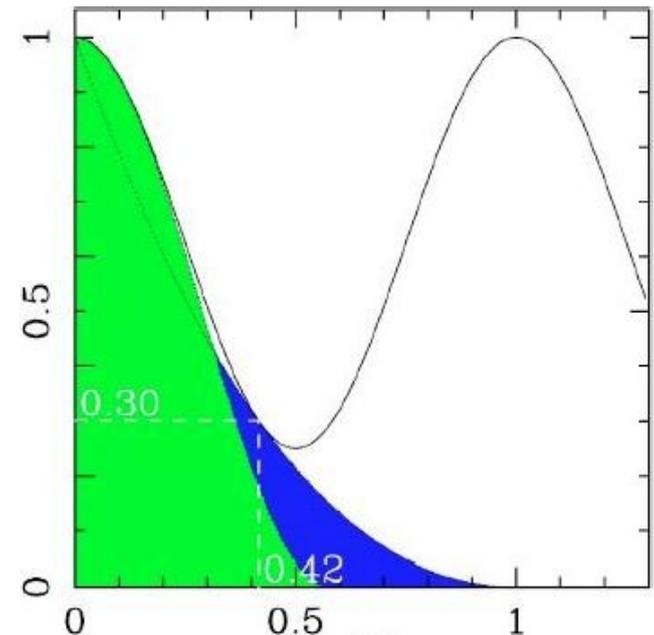
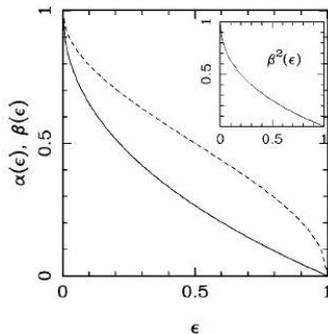
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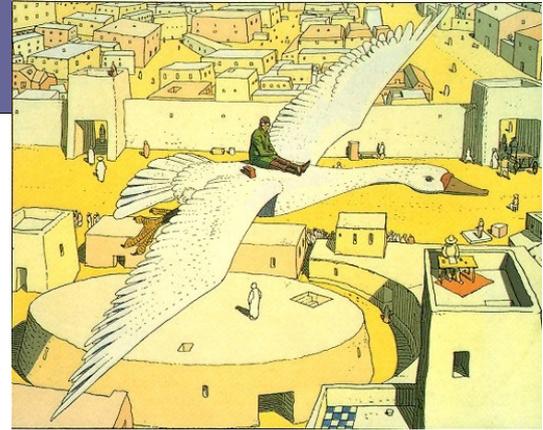
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$$\rightarrow |x' - x| \geq \frac{\pi}{2} \max \left[ \frac{\alpha(\epsilon)}{\mathcal{H}}, \frac{\beta(\epsilon)}{\Delta H} \right]$$

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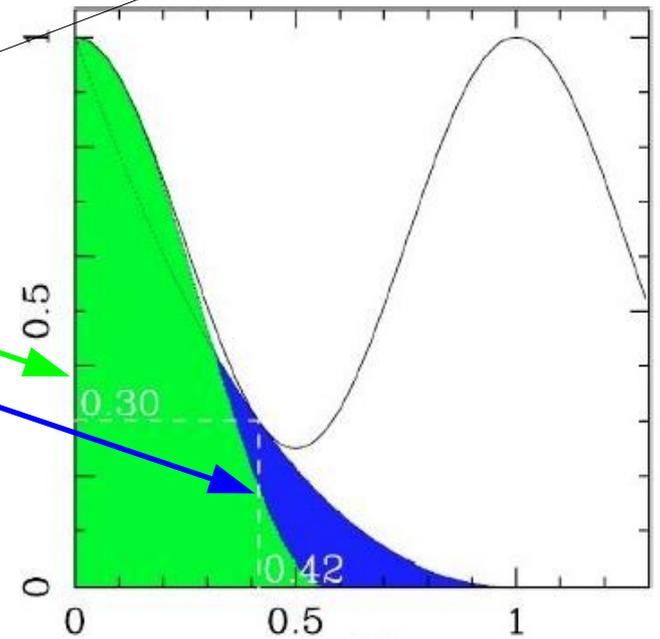
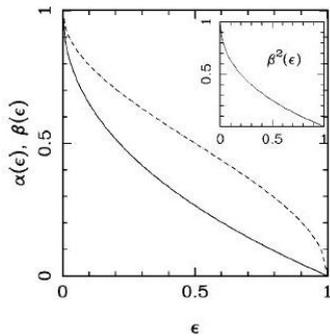
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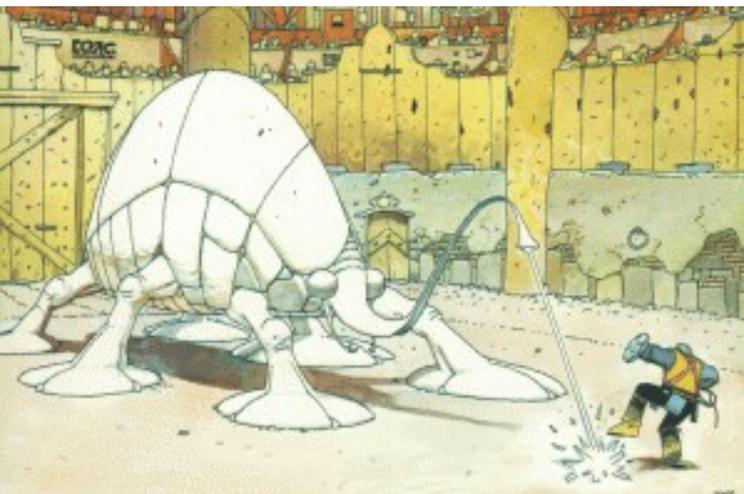


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for  $\nu$  probes,

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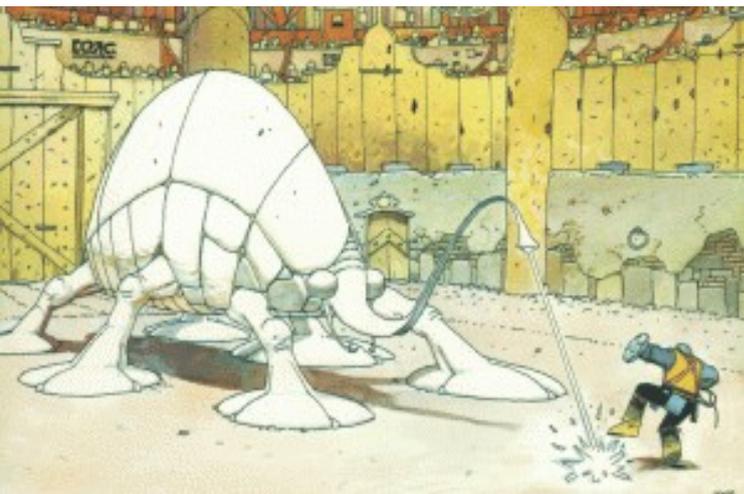
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Cramer-Rao term



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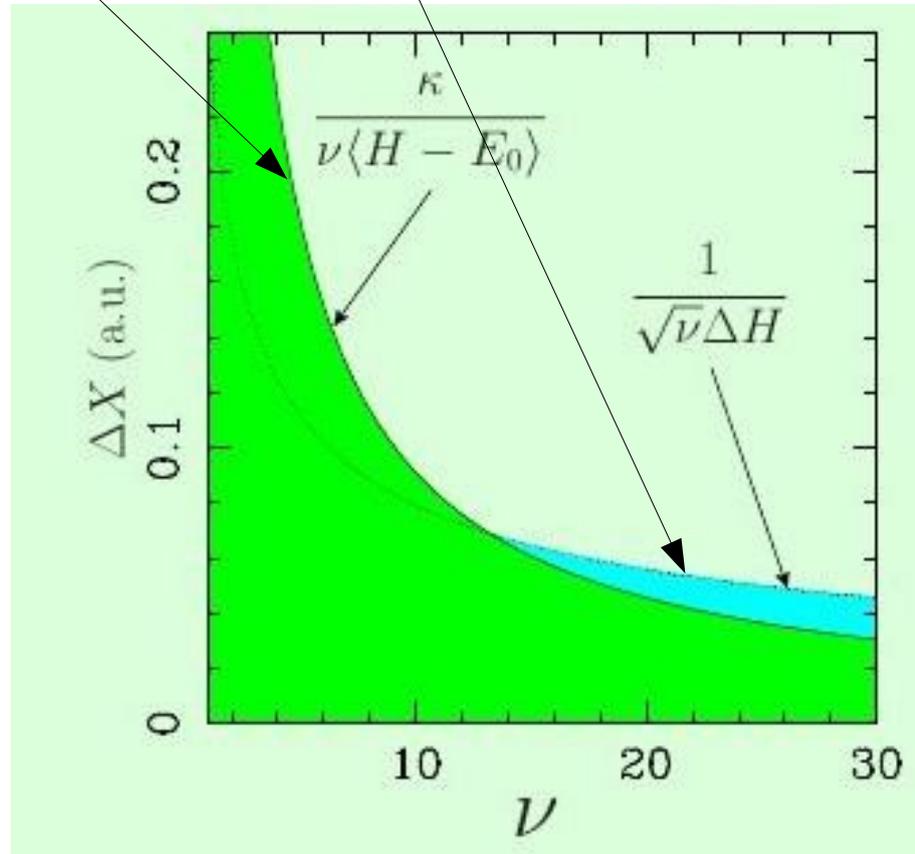
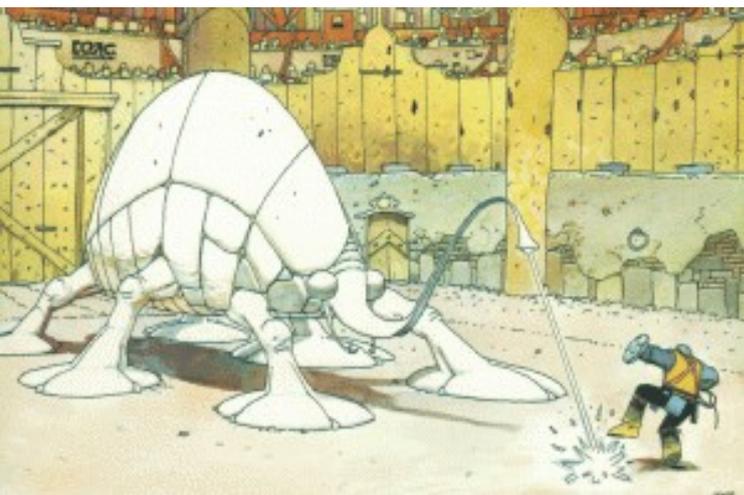
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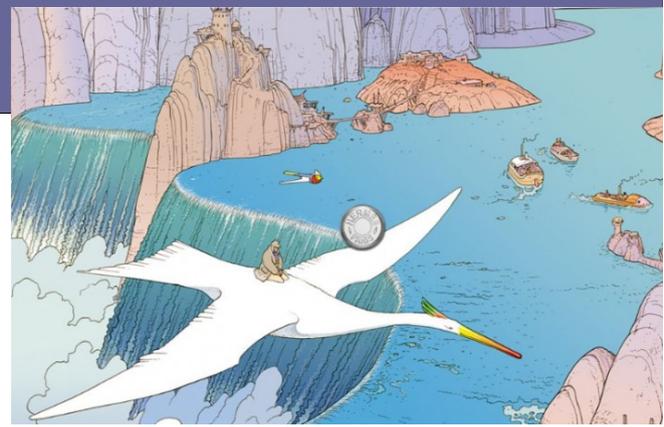
Cramer-Rao term

This is important  $\rightarrow$  Cramer-Rao wins where it is physically significant, for  $\nu \rightarrow \infty$



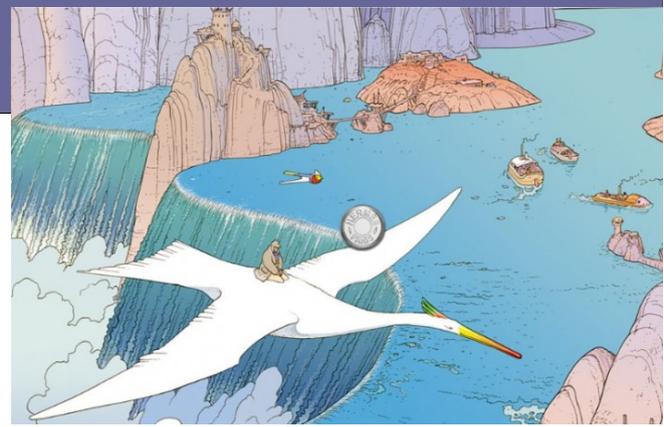
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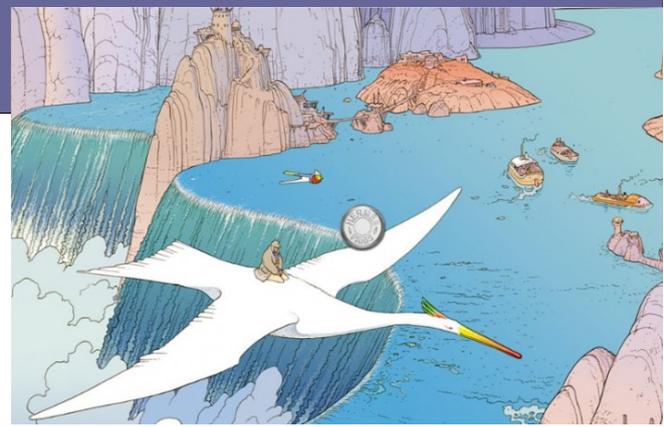


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$$x' - x = 2\lambda \Delta X \Rightarrow F = \epsilon \leq \frac{4}{\lambda^2}$$

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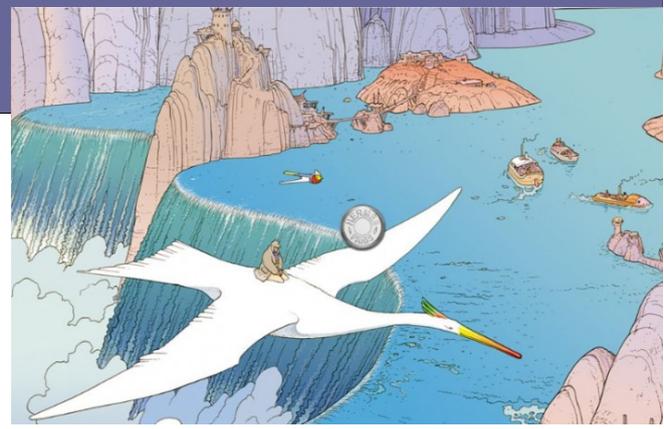


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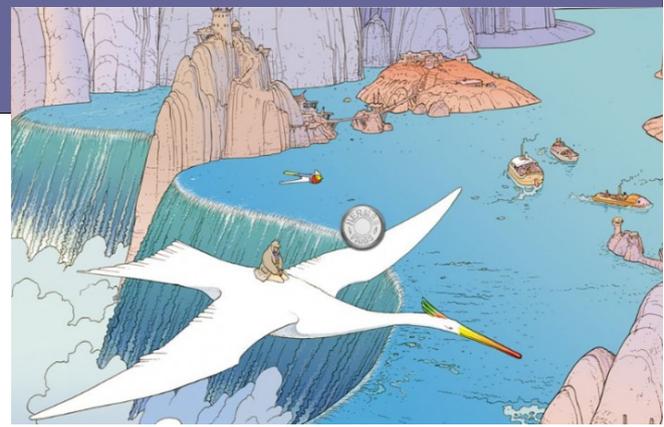
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(and also a Heisenberg-uncertainty bound)

Can we always choose the appropriate  $\lambda$ ?

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e.g., suppose we KNOW the value of  $x$ : there cannot be any bound on the precision of the determination!!! We know it already!

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Our bound fills up **ALMOST** all the space up to the prior information!



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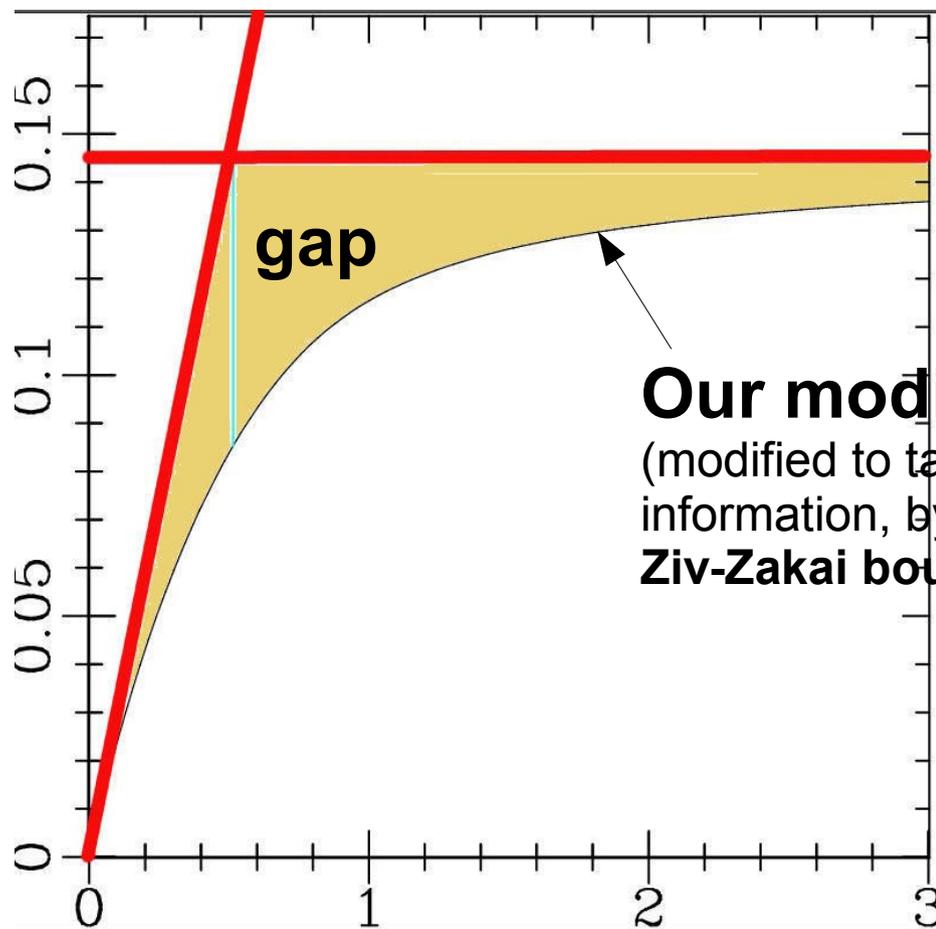
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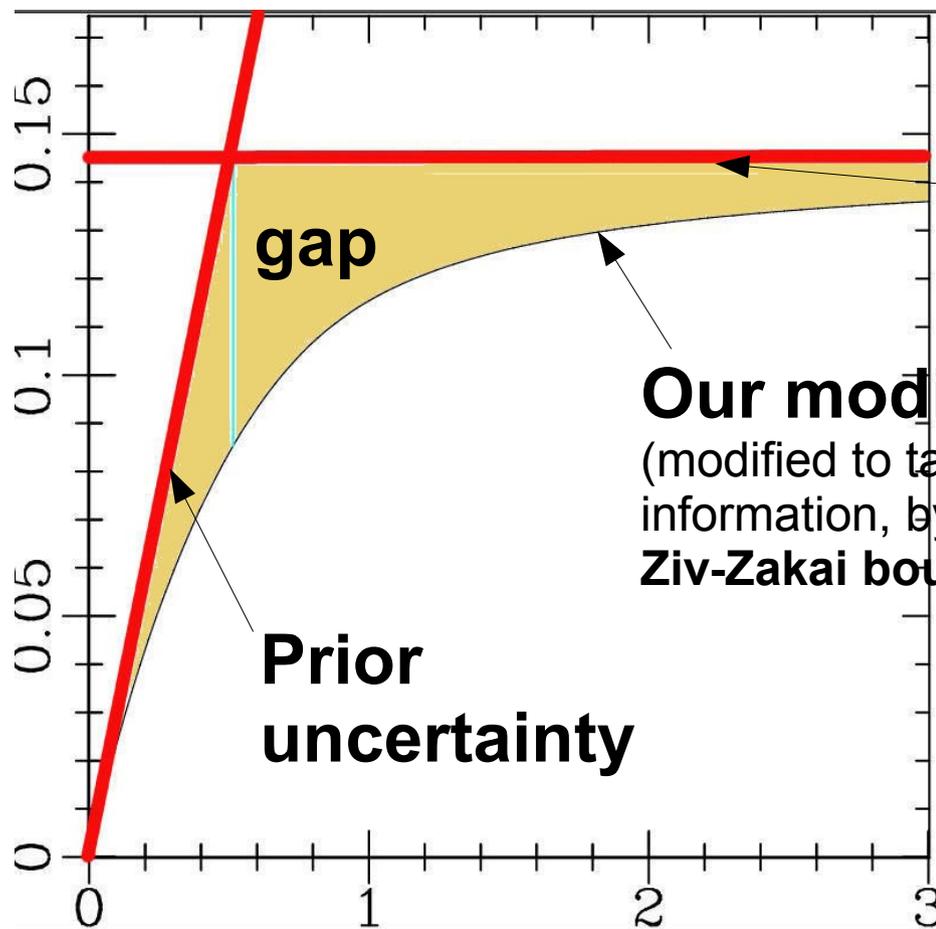
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(modified to take into account the prior information, by joining our bound to the quantum **Ziv-Zakai bound** [Tsang, arXiv:1111.3568])



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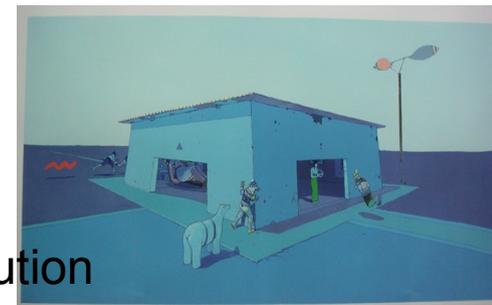
**Heisenberg scaling**

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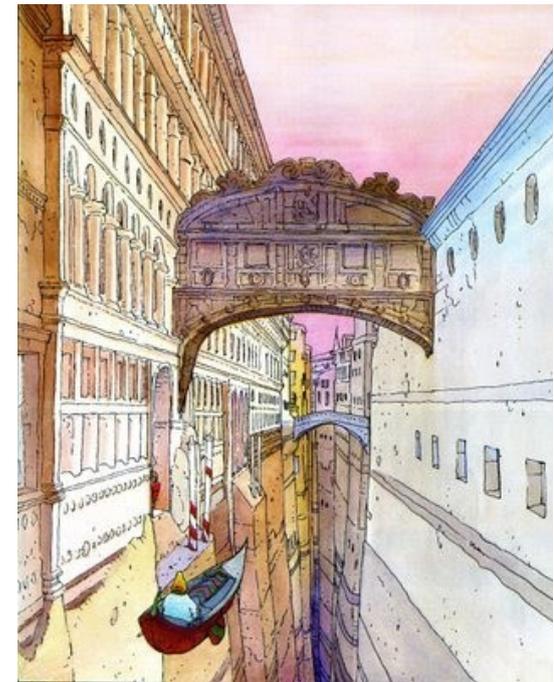
**Prior uncertainty**

w=width of the prior distribution



# Uselessness of sub-Heisenberg estimation

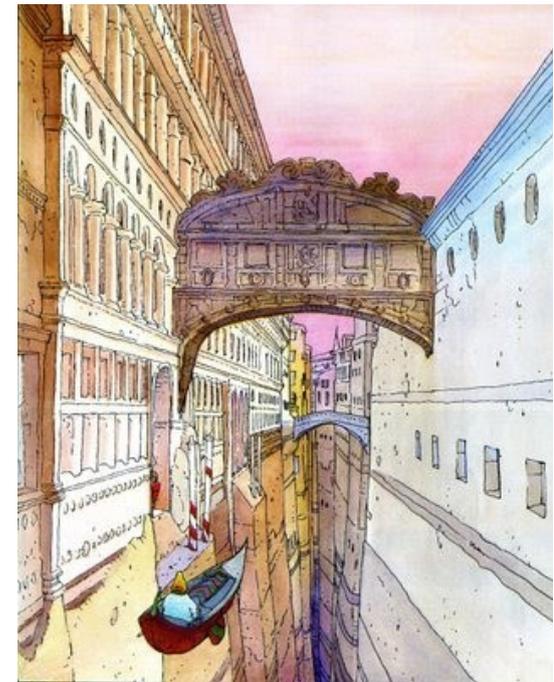
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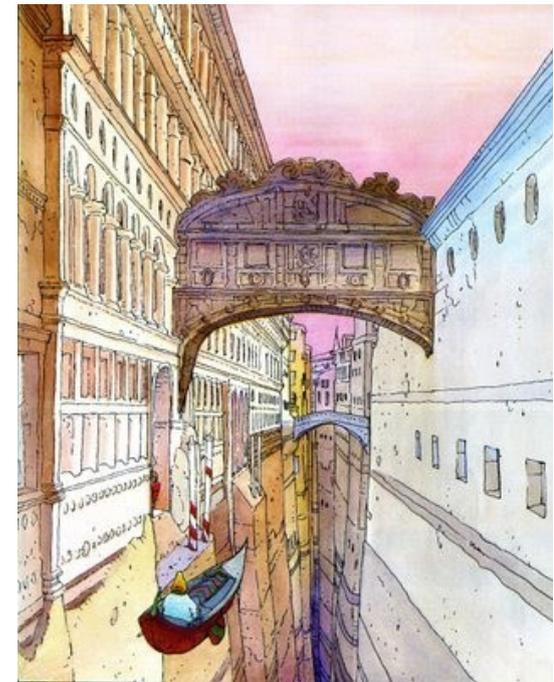


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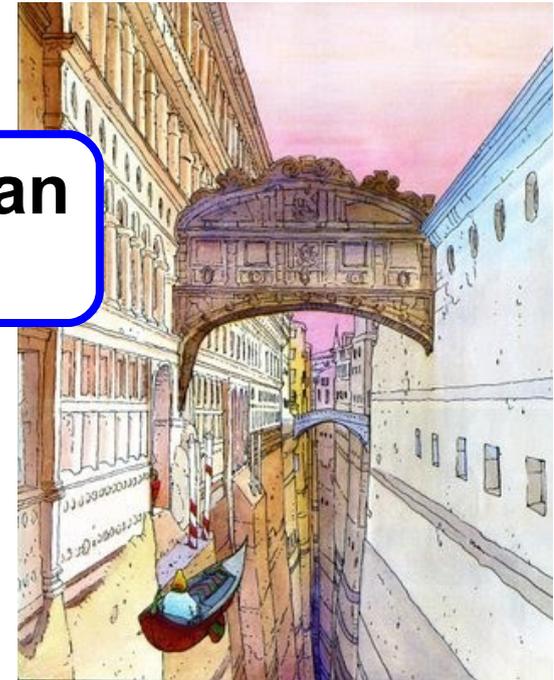
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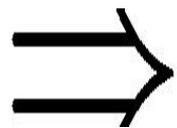
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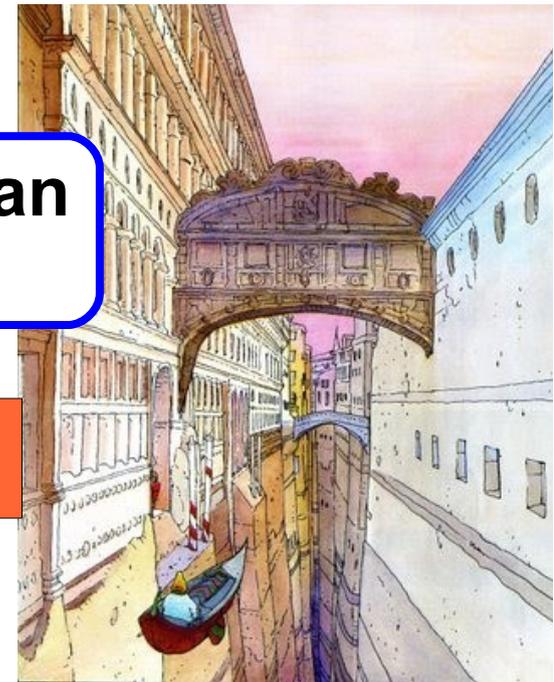
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**sub-Heisenberg strategies are useless**



Incidentally...

**Ziv-Zakai bound**

**+**

**Heisenberg uncertainty/Cramer-Rao part of our bound**



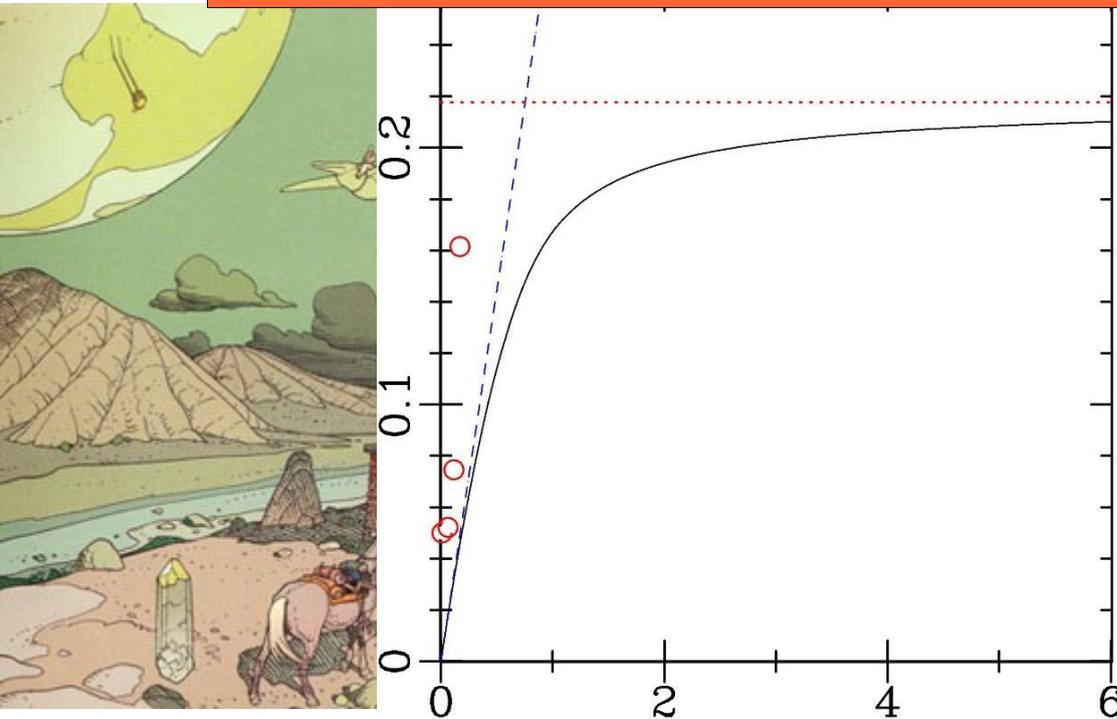
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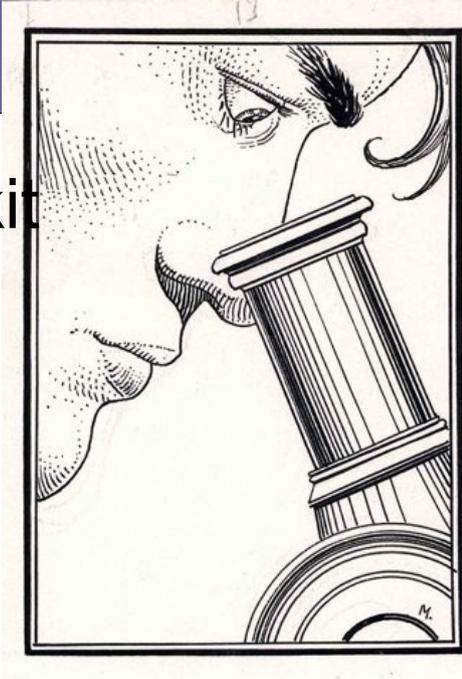
**= an uncertainty relation weighted by an arbitrary prior distribution!**



(Tsang already gave the case of uniform distribution)

# Quantum Metrology

Our bound is part of the **quantum metrology** toolkit



Curious? —▶ Recent review:

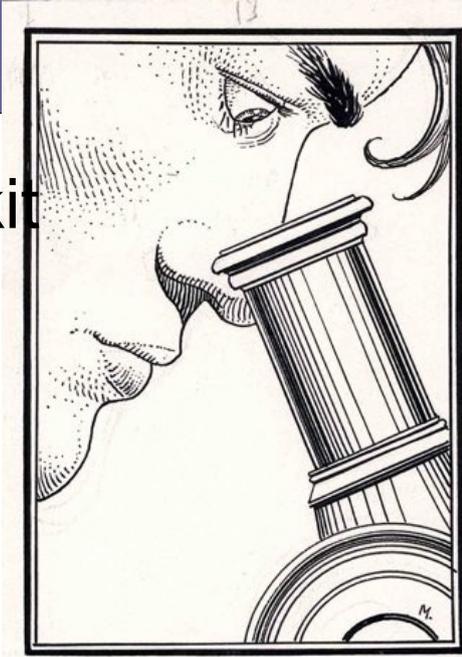
“Advances in Quantum Metrology”, Nature Photonics **5**, 222 (2011).

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quantum technology-oriented

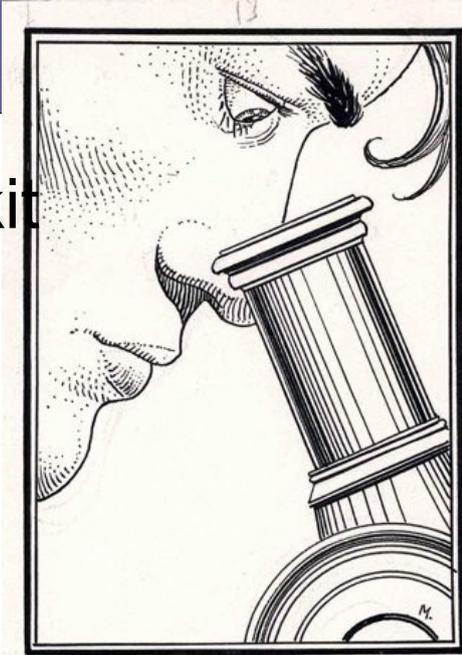


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Determining the ultimate bounds in measurements / measurement problem

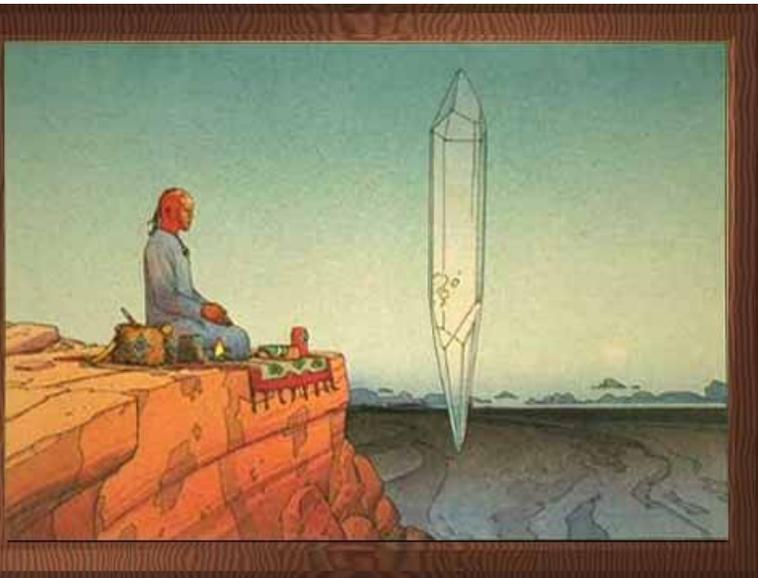
quantum foundations

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# What did I say?!?

- Heisenberg uncertainty relations and q. Cramer-Rao bounds
- Our bound: **a first moment generalization**
- Ideas behind the proof
- Prior information  $\Rightarrow$  sub-Heisenberg strategies exist, but are useless.



A new uncertainty relation  
with EXPECTATION VALUE  
instead of the variance

$$\Delta X (\langle H \rangle - E_0) \geq \kappa$$